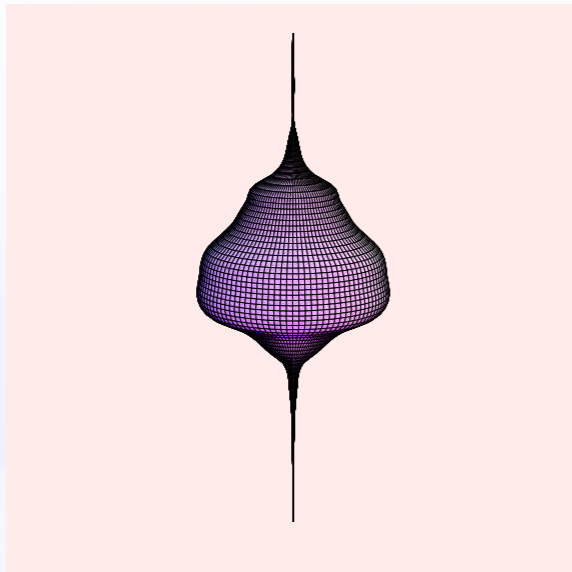




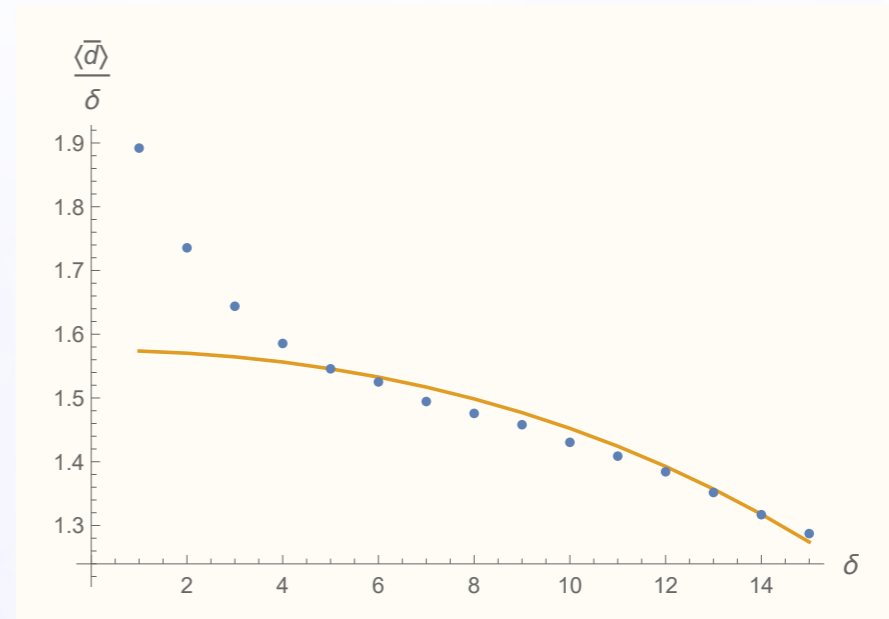
⟨Quantum|Gravity⟩Society

# Quantum Gravity Demystified

Renate Loll



J. Ambjørn, J. Jurkiewicz, R.L., PRL 93 (2004) 131301



N. Klitgaard, R.L., EPJ C 80 (2020) 990

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# My perspective on quantum + gravity

**Aim:** construct a fundamental theory of quantum gravity as a non-perturbative, diffeomorphism-invariant quantum field theory of dynamical geometry and study its properties in a Planckian regime.

This presents major technical, physical and conceptual challenges: dealing with QFT infinities and the absence of a fixed background spacetime, devising appropriate numerical and renormalization methods, (re-)deriving the classical limit and phenomenology.

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***This is possible.*** Major advances towards this goal have been made in the research program of ***Causal Dynamical Triangulations (CDT)***. It sets a concrete frame of reference - beyond taste and style - for what we may expect to be able to achieve in quantum gravity.

# Why should you care?

Talking to various perspectives represented here:

**quantum gravity:** nontrivial, unexpected results despite non-exotic ingredients; functioning computational framework (= *our “lab”*) to evaluate quantum observables beyond perturbation theory; *“CDT is to gravity what lattice QCD is to nonabelian gauge theory”*

**cosmology:** most likely phenomenological predictions will involve early-universe quantum physics, but derived from the full theory *without* an a priori symmetry reduction (unlike quantum cosmology)

**QG experiments in the lab:** will have little to say about weak-gravity regime; lessons on nature of observables, important/subtle to distinguish “physics” and “gauge” (= coordinate effects)

**“fundamental principles”:** quantum (field) theory and general relativity are perfectly compatible; CDT provides bottom-up realization of QG: causal structure is essential, unitarity is realized

# What's the problem with quantum gravity?

- General Relativity = theory *of* spacetime, not *on* (a fixed) spacetime
- quantum theory based on perturbative split  $g_{\mu\nu}(x) = \eta_{\mu\nu}^{\text{Mink}} + h_{\mu\nu}(x)$  on a fixed Minkowskian background is nonrenormalizable

M. Goroff, A. Sagnotti, NPB 266 (1986) 709

- standard relativistic quantum field theory (QFT) not applicable, no blueprint beyond perturbation theory (except nonperturbative lattice QCD, *but* this has a fixed background, different gauge symmetry)
- no experiments or observations to guide theory-building
- (nonperturbative) QG  $\lesssim$  2000: large variety of approaches;  
“**We don't know *what* to compute, and we don't know *how*.”**”
- QG  $\gtrsim$  2000 (post extended-objects era): renaissance of “good old QFT”/the path integral, *we have learned how and what to compute*

R.L. et al.: “Quantum Gravity in 30 Questions”, arXiv: 2206.06762

# Causal Dynamical Triangulations: the basics

- gravitational path integral over metric d.o.f., *nonperturbative* (NP), background-independent, Lorentzian signature, 4D, not “grand-unified”

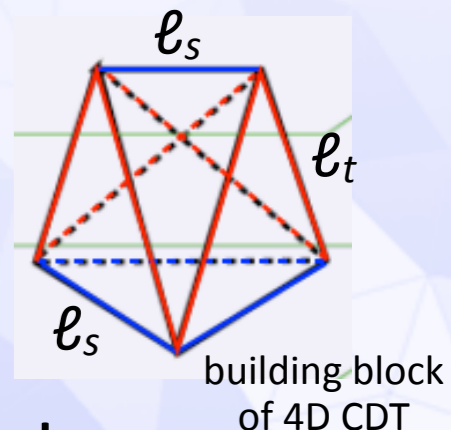
$$Z = \int_{g \in \frac{Lor(M)}{Diff(M)}} \mathcal{D}g e^{iS_{\text{grav}}[g]}$$

- building on Euclidean “dynamical triangulations”, define a new NP Lorentzian 2D path integral: *exactly soluble*  $\Rightarrow$  *signature matters!*

J. Ambjørn, R.L., NPB 536 (1998) 407

- CDT combines emphasis on geometry with path integral covariance (no split  $g_{\mu\nu} = \eta_{\mu\nu} + h_{\mu\nu}$ , no 3 + 1 decomposition)

- uses a regularized version of the space of geometries,  $\mathcal{G}(M) = Lor(M)/Diff(M)$ : piecewise flat, simplicial manifolds  $\mathcal{T}$



- minimal GR ingredients + standard Q(F)T methods, adapted to dynamical geometry + numerical methods = **new territory near  $\mathcal{E}_{PI}$**

- 2D random geometry is hot topic in maths!

S. Sheffield, arXiv:2203.02470

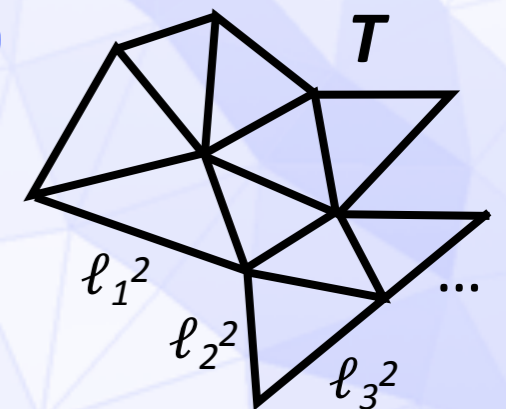
# Putting quantum gravity on a lattice, correctly

**General strategy:** lattice acts as a regulator, with UV cutoff  $a$ ; search for a continuum limit by approaching a second-order phase transition in the limit  $a \rightarrow 0$  while renormalizing bare couplings appropriately; attain *universality* (independence of regularization); this is **not** “discrete QG”

- “reaches where other methods don’t”, subject to numerical limitations; if it exists, continuum theory is essentially *unique*
- “naïve” lattice QG ( $\geq 1979$ ): put various first-order formulations of GR (tetrad  $e_\mu^A$  + spin connection  $\omega_\mu^{AB}$ ) on a fixed hypercubic lattice; problem: diffeomorphism symmetry badly broken; no interesting results
- “not-so-naïve” lattice QG ( $\geq 1981$ ): based on “GR without coordinates”  
 $(M, g_{\mu\nu}(x)) \rightarrow (T, \{\ell_i^2, i=1, \dots, n\}), S_{\text{grav}}[g_{\mu\nu}] \rightarrow S^{\text{Regge}}(T, \{\ell_i^2\})$

T. Regge, *Nuovo Cim.* A19 (1961) 558

- diffeo-invariance manifest, work directly on  $\mathcal{G}(M)$ ;  
CDT ( $\ell^2 = \pm a^2$ ) implementation is labelling-invariant





# The path integral (PI) according to CDT

$$Z = \int_{\mathcal{G}(M)} \mathcal{D}g e^{iS_{\text{grav}}[g]} \rightarrow Z^{\text{CDT}} = \lim_{a \rightarrow 0} \sum_{\substack{\text{inequiv.} \\ \text{causal} \\ \text{triang. } T}} \frac{1}{C(T)} e^{iS^{\text{Regge}}[T]}$$

*bare action*  
# discrete symmetries of T

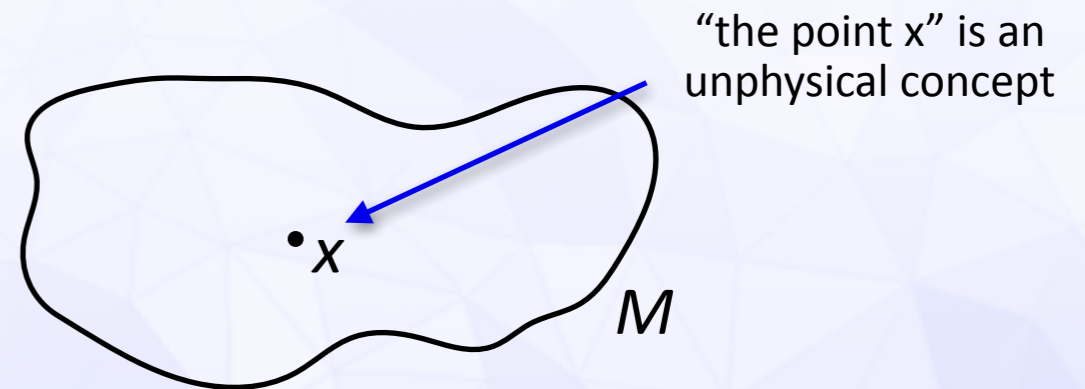
- usually, can't evaluate complex PI, do Euclidean  $\int \mathcal{D}g \exp(-S^{\text{eu}})$  instead
  - ☑ CDT has a well-defined analytic continuation ("Wick-rotation")
- usually, hard to renormalize compatible with diffeomorphism symmetry
  - ☑ CDT has no residual symmetries, has a geometric cutoff  $a$
- usually, PI highly divergent, no unique renormalization;
  - ☑ numerical evidence of exponential bound on # of configurations
- usually, cannot do any computations, PI not Gaussian
  - ☑ CDT amenable to Monte Carlo simulation; get quantitative results
- usual problem: why should PI lead to a unitary theory?
  - ☑ CDT reflection-positive w.r.t. discrete "proper time"

# CDT quantum gravity: results

- we have a computational framework — what can we do with it?
- physics of **quantum spacetime** is captured by invariant **quantum observables**  $\hat{\mathcal{O}}$ :

$$\langle \hat{\mathcal{O}} \rangle = \frac{1}{Z} \int \mathcal{D}g \mathcal{O}[g] e^{-S_{\text{grav}}[g]}$$

- observables in Yang-Mills theory are local scalars, like  $F^{\mu\nu}F_{\mu\nu}$ , but observables in pure gravity are nonlocal integrals of scalars, like  $\int_M d^4x \sqrt{g} R(x)$



- “expectation management”: your favourite (semi-)classical question will not have a Planckian implementation (this is a feature)
- quantum gravity signature: CDT predicts a reduction  $4 \rightarrow 2$  of the **spectral dimension** @  $\ell_{Pl}$ , [J. Ambjørn, J. Jurkiewicz, R.L., PRL 95 \(2005\) 171301](#) reproduced across approaches — universal in QG? [S. Carlip, CQG 34 \(2017\) 193001](#)

# Key result: emergence of classicality from CDT

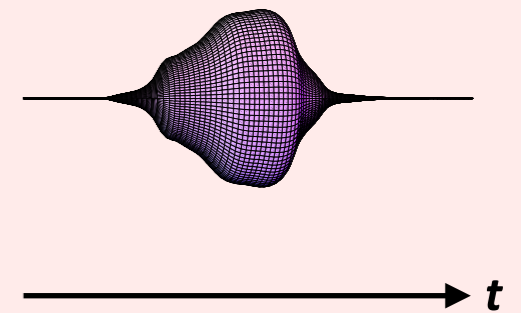
The measured average shape  $\langle V_3(t) \rangle$  (spatial volume as a function of proper time) of the dynamically generated quantum spacetime in CDT matches that of a classical de Sitter space.

J. Ambjørn, A. Görlich, J. Jurkiewicz, R.L., PRL 100 (2008) 091304, PRD 78 (2008) 063544

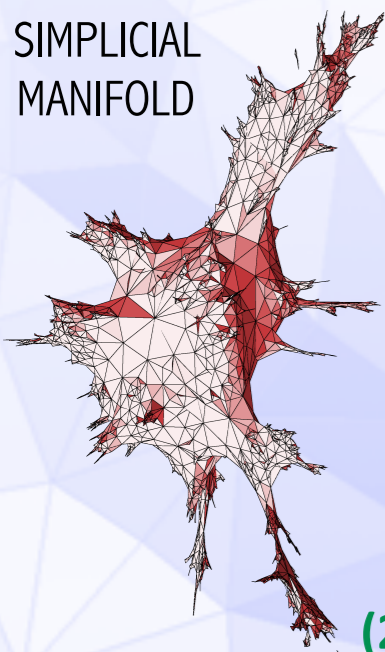
Since the global shape of the universe is just a single mode of the metric, we cannot conclude that it *is* a (Euclidean) de Sitter space  $S^4$ , with line element

$$ds^2 = dt^2 + c^2 \cos^2(t/c) d\Omega_{(3)}^2.$$

MC snapshot of the shape  $\langle V(t) \rangle$  of the universe



SIMPLICIAL  
MANIFOLD

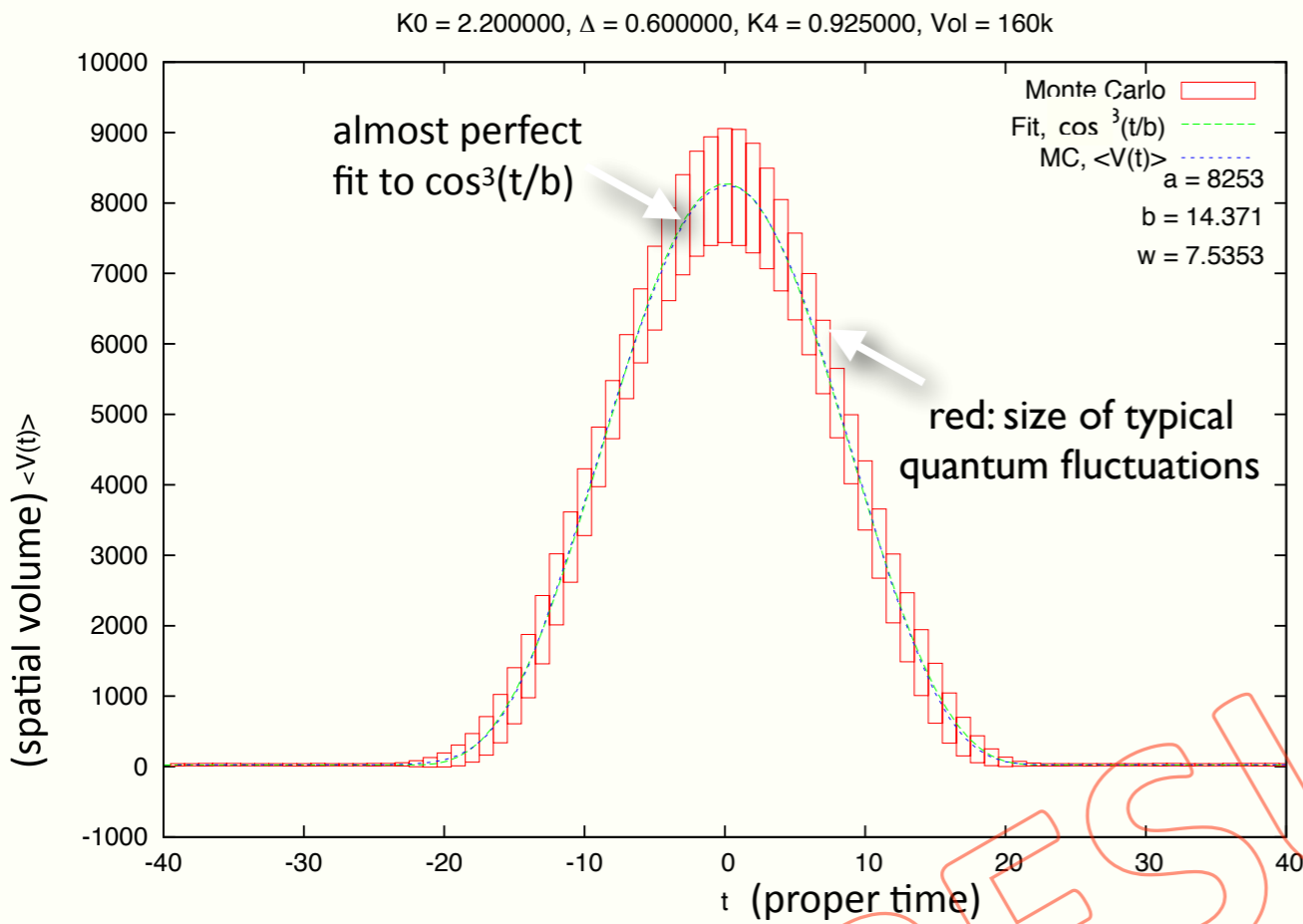


What about the **local geometry** of this quantum universe?

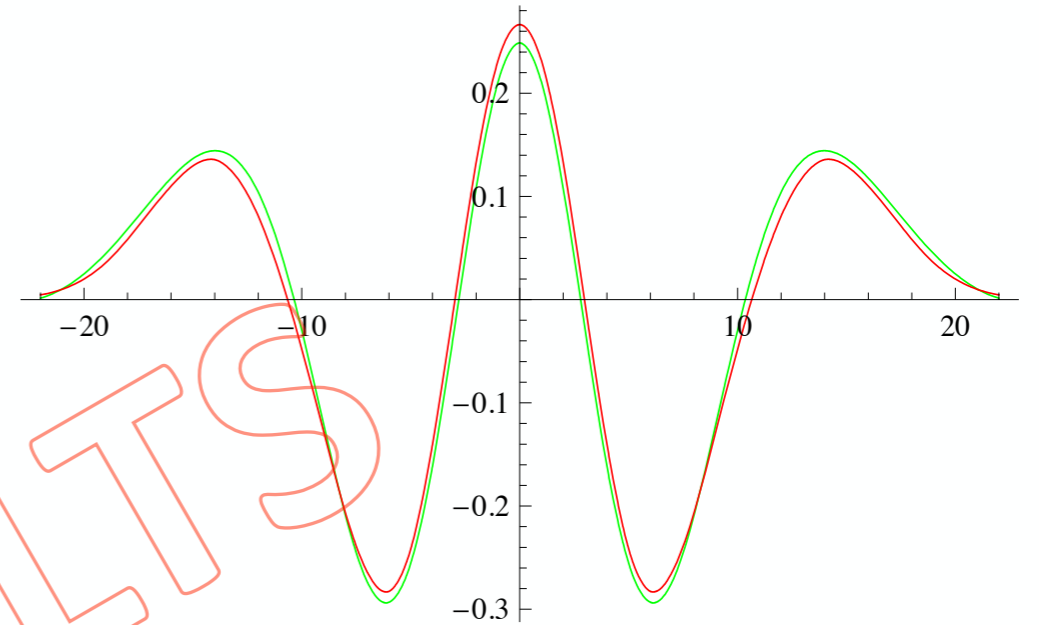
Can we attribute **local curvature** to a non-smooth metric space?  $R^{\kappa}_{\lambda\mu\nu}(x) = [g, \partial g, \partial^2 g] = ?$  Recently, we defined, tested and measured a well-defined notion of **quantum Ricci curvature** applicable in a Planckian regime. N. Klitgaard & RL, PRD 97

(2018) 0460008 and 106017, Eur. Phys. J. C80 (2020) 990, J. Brunekreef & RL, PRD 103 (2021) 026019

# The universe is de Sitter-shaped

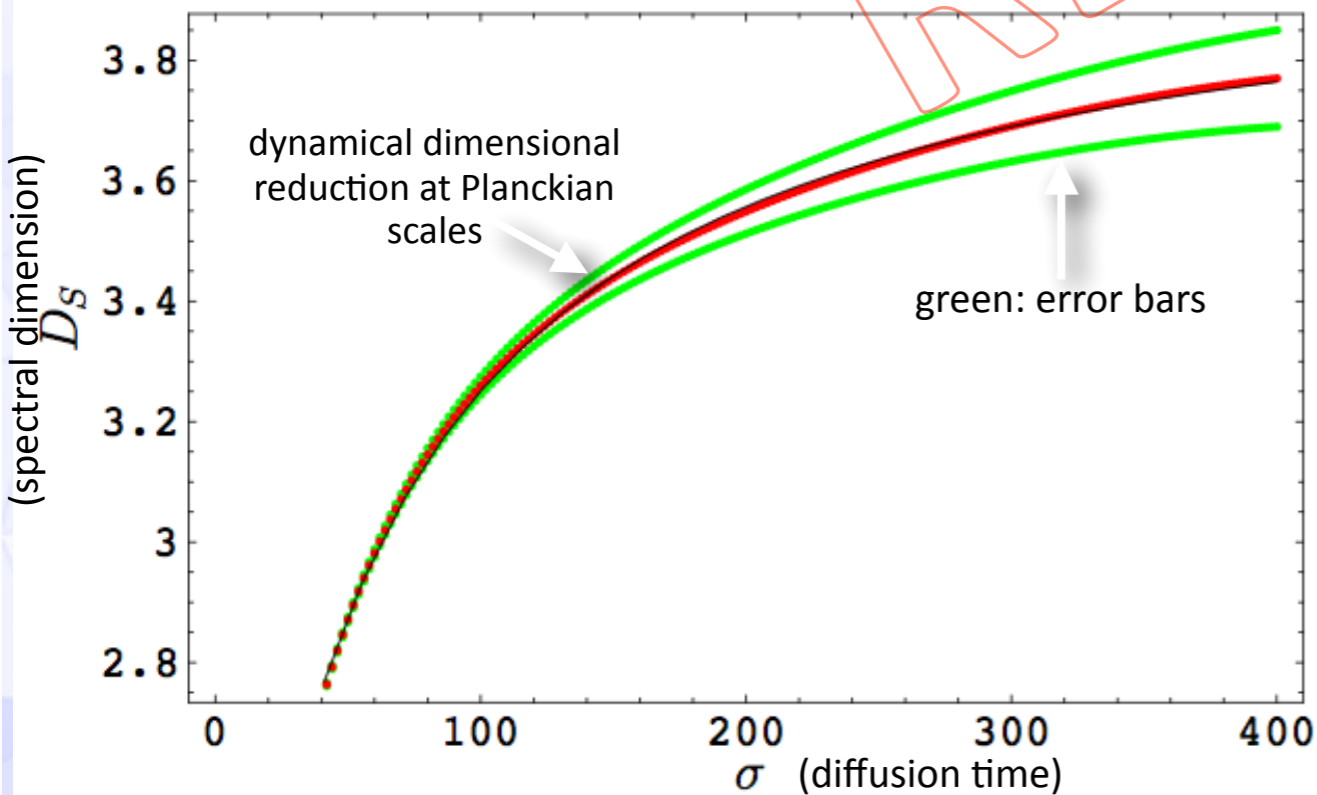


# Volume fluctuations around de Sitter

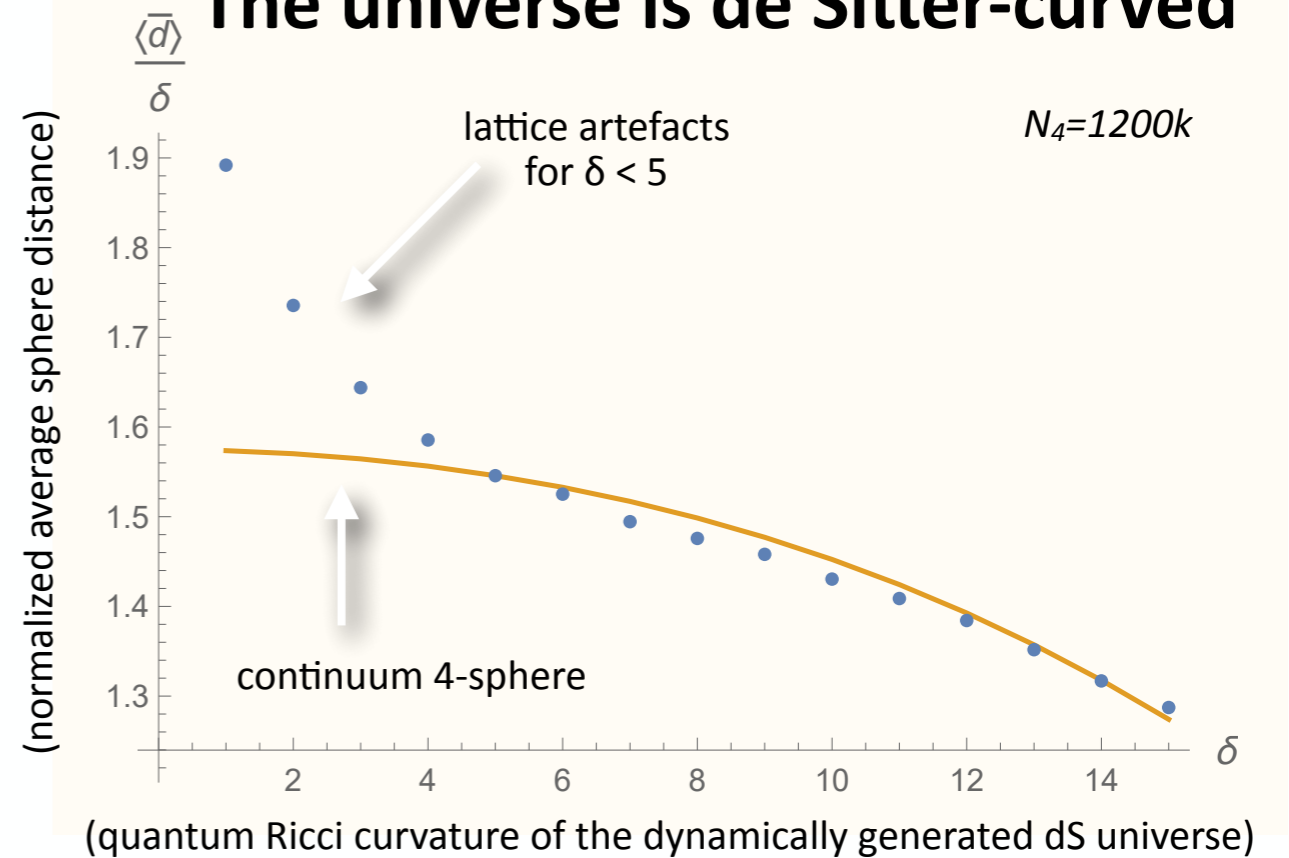


(low-lying eigenmode matches with semiclassical expectation)

# Spectral dimension of the universe



# The universe is de Sitter-curved



# Relation to our actual universe

CDT predicts a universe **with positive  $\Lambda$** , which on large scales is extended and four-dimensional, and whose **shape and average curvature** are compatible with those **of a de Sitter space**, matching our current understanding of the very early universe.

Remarkably, these properties have been derived *from first principles* in the full quantum theory; we also have in principle access to (diffeomorphism-invariant) correlation functions.

At what scales and how does gravity interact with matter?

$$Z = \int_{\mathcal{G}(M)} \mathcal{D}g \int_{\Phi} \mathcal{D}\phi e^{i(S_{\text{grav}}[g] + S_{\text{matter}}[g, \phi])}$$

Investigations of CDT coupled to matter fields have not found a significant impact on the geometry  $\Rightarrow$  “matter doesn’t matter at  $\ell_{Pl}$ ”?

# Current ambitions and prospects

CDT QG is in a position to reap the benefits of a nonperturbative framework that can produce “numbers” (= expectation values of quantum observables) without relying on ad-hoc assumptions. The art is to ***identify (more) observables*** that can be reliably measured inside the available scale window, while yielding interesting physics.

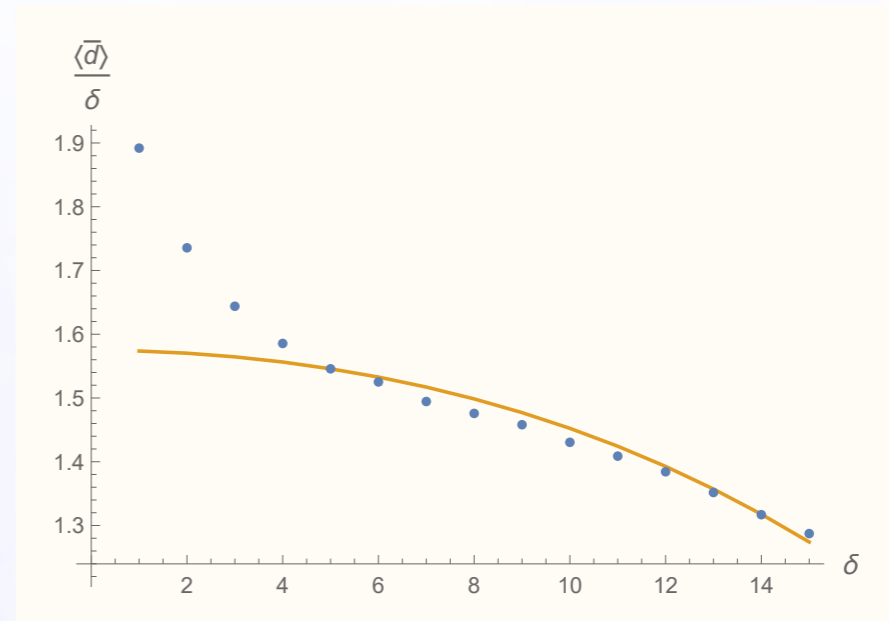
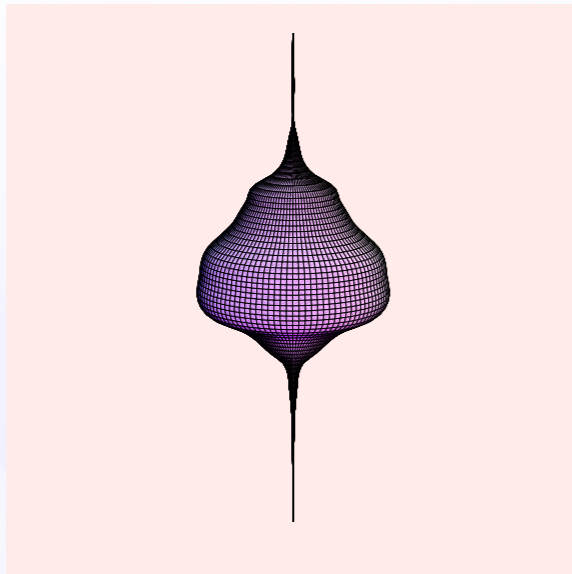
The new ***quantum Ricci curvature*** opens exciting avenues towards a relation with early-universe physics:

- use it to quantify the *local* effect of a massive Planckian particle and compare with a semiclassical expectation [G. Clemente, R.L., w.i.p.](#)
- use it to examine the string-like singularity spontaneously forming in the bifurcation phase of CDT, as a possible *candidate of early-universe structure formation (a primordial black hole?)*
- compute its *two-point function* and compare with QFT on de Sitter space [J. van der Duin, R.L., w.i.p.](#)

# Outlook

- ***genuine progress*** in nonperturbative quantum gravity: instead of comparing “approaches”, started to compare observables/results, e.g. with functional RG methods [F. Saueressig & collaborators](#)
  - CDT quantum gravity is a rare example of ***spacetime emergence***
  - ***work in progress***: quantum measures of homogeneity & isotropy [A. Silva, R.L.](#); extend earlier RG flow analysis [J. Ambjørn, J. Gizbert-Studnicki, A. Görlich, J. Jurkiewicz, R.L., Front. in Phys. 8 \(2020\) 247](#) and look for independent evidence of asymptotic safety; measure curvature and its correlators on spatial slices [J. Brunekreef, R.L., to appear, ...](#)
  - ***challenge***: match nonperturbative and perturbative observables
- ⇒ watch this space!

**CDT reviews**: [J. Ambjørn, A. Görlich, J. Jurkiewicz, R.L., Phys. Rep. 519 \(2012\) 127, arXiv: 1203.3591](#); [R.L., Class. Quant. Grav. 37 \(2020\) 013002, arXiv:1905.08669](#)



**Thank you!**





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